

### Outline for Day 13 - 14

Office hours: 2 - 3

In which we solve the Schrödinger Equation for the hydrogen atom.

- Bohr atom + de Broglie postulate
- Hydrogen atom – a 3 dimensional spherically symmetric object
  - Spherical coordinates
  - Schrodinger equation in 3 dimensions and spherical coordinates
- Solving the hydrogen atom
  - Solution for  $g(\phi)$
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  - Solution for  $f(\theta)$
  - Complete solutions

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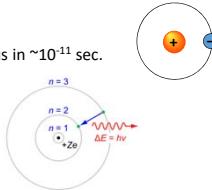
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### Models of the Atom



- Thomson – “Plum Pudding”
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  - Problem: Doesn't match scattering experiments
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  - Why? Scattering showed a small, hard core.
  - Problem: electrons should spiral into nucleus in  $\sim 10^{-11}$  sec.
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  - Why? Explains spectral lines.
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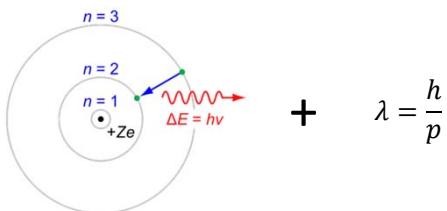
### Successes of Bohr Model

- 'Explains' source of Balmer formula and predicts empirical constant  $R$  (Rydberg constant) from fundamental constants
- Explains why  $R$  is different for different single electron atoms (called *hydrogen-like ions*).
- Predicts approximate size of hydrogen atom
- Explains (sort of) why atoms emit discrete spectral lines
- Explains (sort of) why electron doesn't spiral into nucleus

### Shortcomings of the Bohr model:

- Why is angular momentum quantized yet Newton's laws still work?
- Why don't electrons radiate when they are in fixed orbitals yet Coulomb's law still works?
- No way to know *a priori* which rules to keep and which to throw out...
- Can't explain shapes of molecular orbitals and how bonds work
- Can't explain doublet spectral lines

### Bohr Atom + de Broglie electron waves



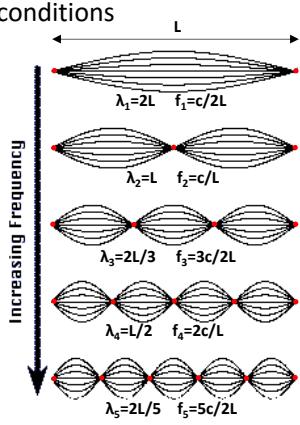
$$\lambda = \frac{h}{p}$$

### Waves with boundary conditions

- Boundary conditions mean that waves only have discrete frequencies.
- e.g., guitar strings

$$\lambda_n = 2L/n \quad f_n = nc/2L$$

• node = fixed point that doesn't move

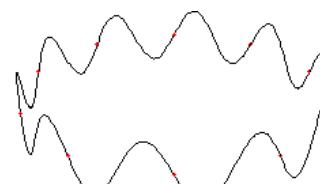


PHET: normal modes

### Standing Waves on a Ring

Just like standing wave on a string,  
but now the two ends of the string are joined.

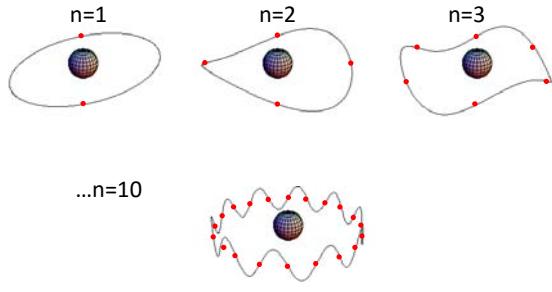
$$2\pi r = n\lambda \quad n = 1, 2, 3, \dots$$



Derive angular momentum of an electron wave

deBroglie electron waves in an atom

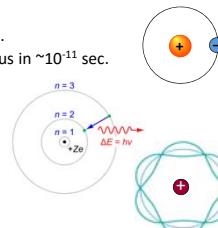
$$\lambda = \frac{h}{mv}$$



• = node = fixed point that doesn't move

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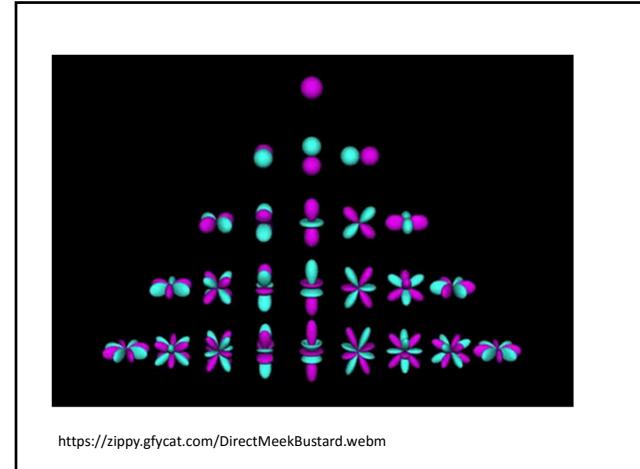
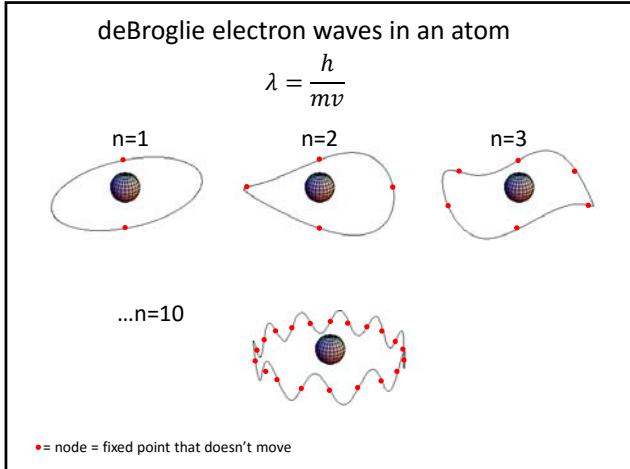


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- deBroglie – electron standing waves
  - Why? Explains quantized energy levels
  - Problem: still only works for Hydrogen.
- Schrodinger – quantum wave functions
  - Why? Explains everything!
  - Problem: None (except that it’s abstract)

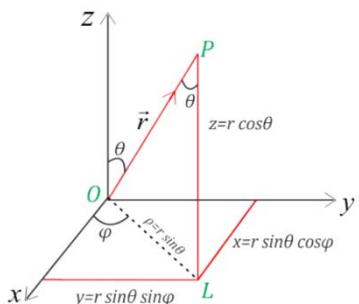
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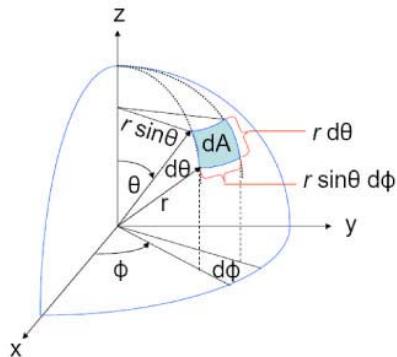
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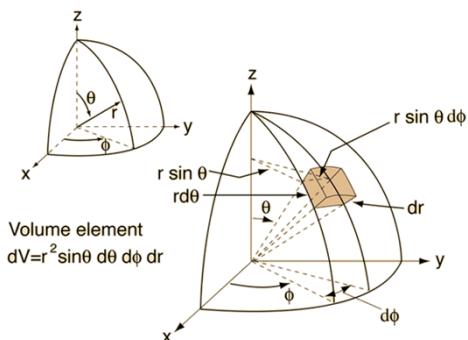
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Schrödinger equation in 3D spherical coordinates

$$\frac{-\hbar^2}{2m} \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial \psi}{\partial r} - \frac{\hbar^2}{2mr^2} \left[ \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \right] + U\psi = E\psi \quad \text{Eqn. 8.47}$$

prove that:  $\frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial \psi}{\partial r} = \frac{1}{r} \frac{\partial^2}{\partial r^2} r\psi$

$$\frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial \psi}{\partial r} = \frac{1}{r^2} \left[ 2r \frac{\partial \psi}{\partial r} + r^2 \frac{\partial^2 \psi}{\partial r^2} \right] = \frac{2}{r} \frac{\partial \psi}{\partial r} + \frac{\partial^2 \psi}{\partial r^2}$$

$$\frac{1}{r} \frac{\partial^2}{\partial r^2} r\psi = \frac{1}{r} \frac{\partial}{\partial r} \left( \psi + r \frac{\partial \psi}{\partial r} \right) = \frac{1}{r} \left( \frac{\partial \psi}{\partial r} + \frac{\partial \psi}{\partial r} + r \frac{\partial^2 \psi}{\partial r^2} \right) = \frac{2}{r} \frac{\partial \psi}{\partial r} + \frac{\partial^2 \psi}{\partial r^2}$$

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Maybe a solution could look like

$$\psi = R(r)f(\theta)g(\phi)$$

$$-\frac{\hbar^2}{2m} \frac{f g}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} R(r) - \frac{\hbar^2}{2mr^2} \left[ \frac{R_g}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial f}{\partial \theta} + \frac{R_f}{\sin^2 \theta} \frac{\partial^2 g}{\partial \phi^2} \right]$$

$$-\frac{k^2 e^2}{r} R f g = E R f g$$

multiply both sides by  $-\frac{2mr^2 \sin^2 \theta}{\hbar^2 R f g}$

if you can show that  $f(x) = g(y)$  for all  $x$  and  $y$   
 $f(x) = g(y) = \text{constant}$

$$\frac{\sin^2 \theta}{R} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} R + \frac{\sin \theta}{f} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial f}{\partial \theta} + \frac{1}{g} \frac{\partial^2}{\partial \phi^2} g = -\frac{2mr^2 E \sin^2 \theta}{\hbar^2}$$

$$\frac{1}{g} \frac{\partial^2 g}{\partial \phi^2} = -\frac{\sin^2 \theta}{R} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} R - \frac{\sin \theta}{f} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial f}{\partial \theta} - \frac{2mr^2 \sin^2 \theta}{\hbar^2} (E + \frac{k^2 e^2}{r})$$

is a function of  $\phi$       is a function of  $R$  and  $\theta$

$$\frac{1}{g} \frac{\partial^2 g}{\partial \phi^2} = \text{constant} = -m^2 \Rightarrow \frac{\partial^2 g}{\partial \phi^2} = -m^2 g \quad \text{let } g = e^{im\phi}$$

then  $\frac{\partial g}{\partial \phi} = im e^{im\phi} \quad \frac{\partial^2 g}{\partial \phi^2} = (im)^2 e^{im\phi} = -m^2 g$

boundary condition for  $\phi$ :  $g(\phi + 2\pi) = g(\phi)$   
 Single valued

$$e^{im\phi} = e^{im(\phi + 2\pi)} = e^{im\phi} e^{im2\pi} \quad m=0, \pm 1, \pm 2, \dots$$

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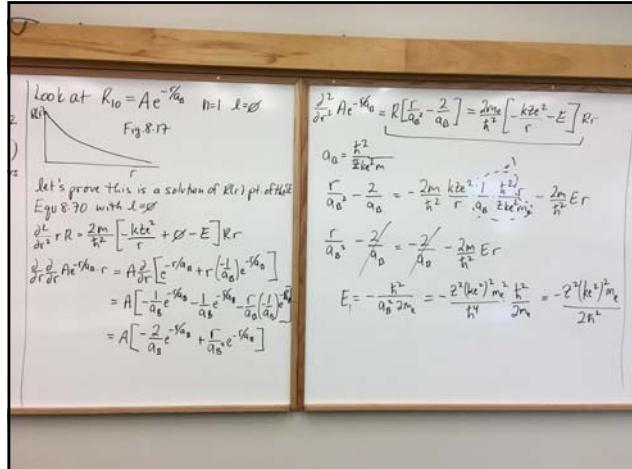
$$\begin{aligned}
 & \text{Move all } r \text{ dependencies to one side} \\
 & \frac{1}{R} \frac{\partial}{\partial r} r^2 \frac{\partial R}{\partial r} + \frac{2mr^2}{\hbar^2} \left( E + \frac{kze^2}{r} \right) = \\
 & \quad \text{depends only on } r \\
 & \text{Multiply both sides by } \frac{R}{\hbar^2} \\
 & \text{*use the identity for radial derivative term} \\
 & \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial R}{\partial r} = l(l+1) \frac{R}{\hbar^2} \frac{2m}{2M} \frac{z^2}{r^2} - \\
 & \frac{\partial^2}{\partial r^2} r^2 R = \frac{2m}{\hbar^2} \left[ -\frac{kze^2}{r} + \frac{l(l+1)\hbar^2}{2mr^2} \right] R r \quad \text{Equ 8.70}
 \end{aligned}$$

Equ. 8.70 is a differential equation that has solutions tabulated in Table 8.2  
 boundary conditions (not shown) require:  $l < n$   $l$  and  $n$  integers  
 so  $n=1 \Rightarrow l=0$   
 $n=2 \Rightarrow l=0, 1$   
 that the energy is quantized  
 $E_n = -\frac{m_e (ke^2)^2}{2\hbar^2} \frac{1}{n^2} = -13.6 \text{ eV} \frac{1}{n^2}$  Equ 8.74

 $R_{nl}(r)$ : The Radial Wave Functions for the atom

TABLE 8.2  
 The first few radial functions  $R_{nl}(r)$  for the hydrogen atom. The variable  $\rho$  is an abbreviation for  $\rho = r/a_B$  and  $a$  stands for  $a_B$ .

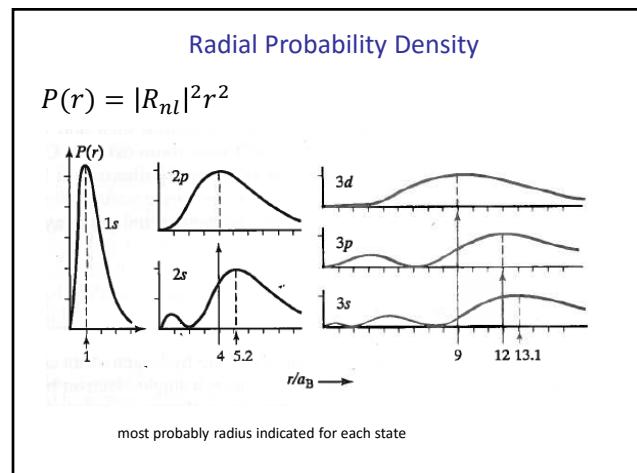
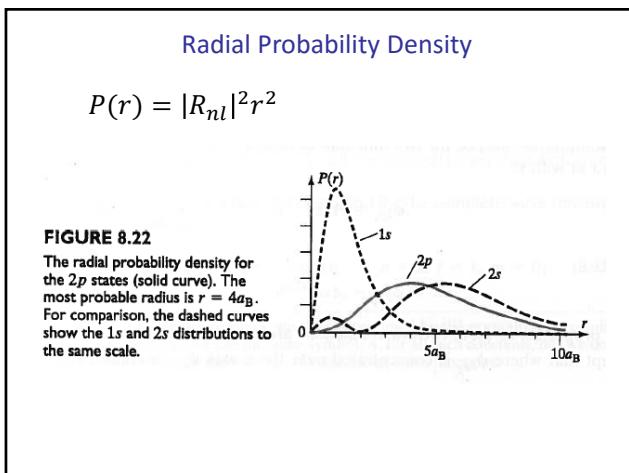
	$n = 1$	$n = 2$	$n = 3$
$l = 0$	$\frac{2}{\sqrt{a^3}} e^{-\rho}$	$\frac{1}{\sqrt{2a^3}} \left( 1 - \frac{1}{2}\rho \right) e^{-\rho/2}$	$\frac{2}{\sqrt{27a^3}} \left( 1 - \frac{2}{3}\rho + \frac{2}{27}\rho^2 \right) e^{-\rho/3}$
$l = 1$		$\frac{1}{\sqrt{24a^3}} \rho e^{-\rho/2}$	$\frac{8}{27\sqrt{6a^3}} \left( 1 - \frac{1}{6}\rho \right) \rho e^{-\rho/3}$
$l = 2$			$\frac{4}{81\sqrt{30a^3}} \rho^2 e^{-\rho/3}$

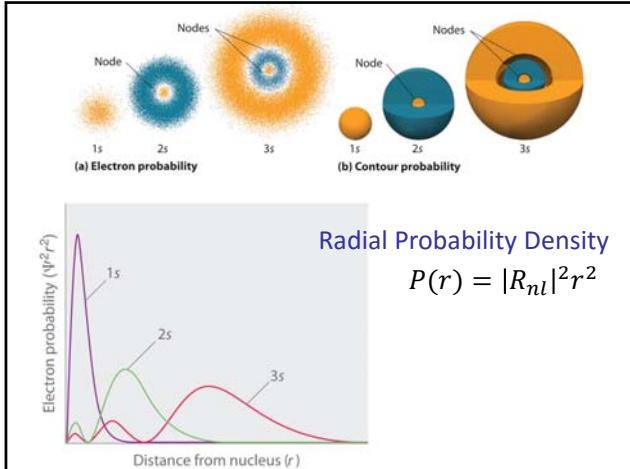


$$P(r) = |R_{nl}|^2 dV = |R_n|^2 r^2 \sin \theta d\theta d\phi$$

$$P(r) = |\psi|^2 dr$$

$$|\psi|^2 dr$$





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Move all  $r$  dependence to one side

$$\frac{1}{R} \frac{d}{dr} r^2 \frac{dR}{dr} + \frac{2mr^2}{\hbar^2} \left( E + k_2 e^2 \right) =$$

depends only on  $r$

$$l(l+1) =$$

Solutions are in table 8.1  
boundary conditions restrict values of  $m$  and  $l$  such that:  
 $l \geq 0$     $|m| \leq l$     $m$ ,  $l$  integers

Often angular functions are written together

$$Y_{lm}(\theta, \phi) = f_{lm}(\theta) g_m(\phi)$$

$$= f_{lm}(\theta) e^{im\phi}$$

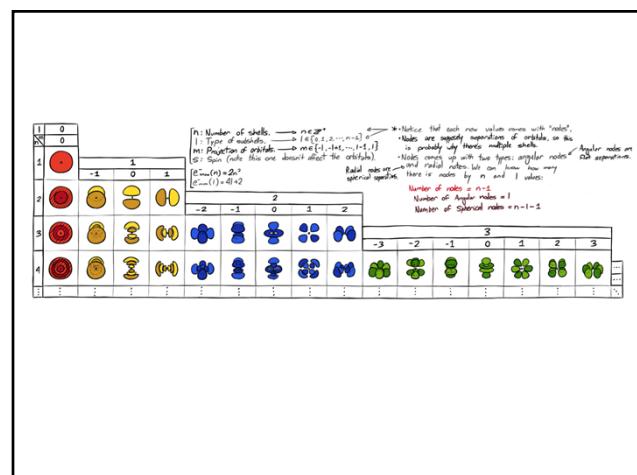
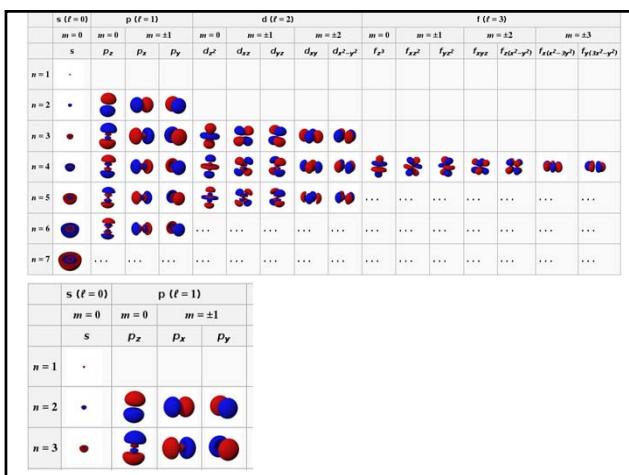
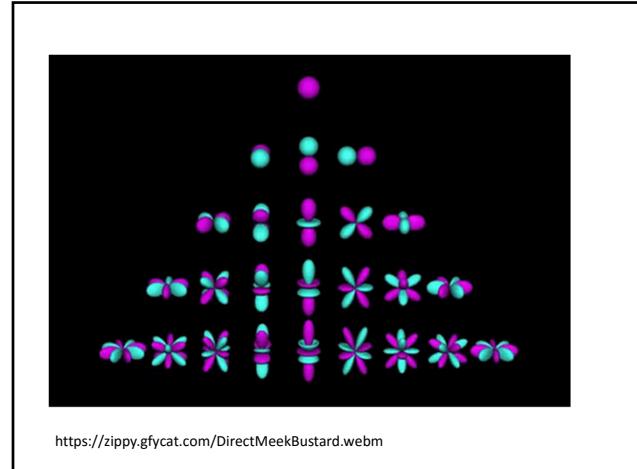
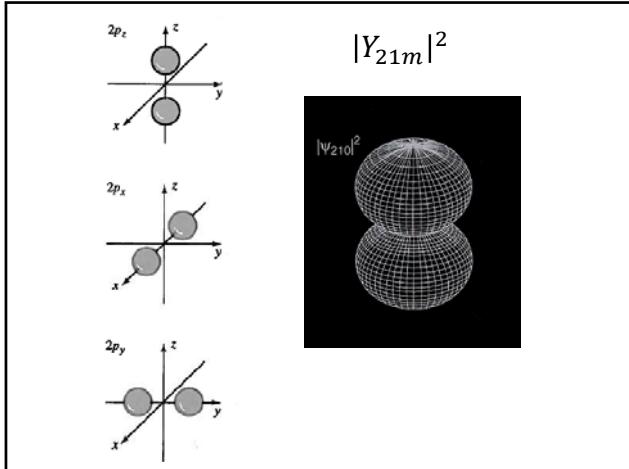
### $Y_{lm}(\theta)$ : The spherical harmonic wave functions

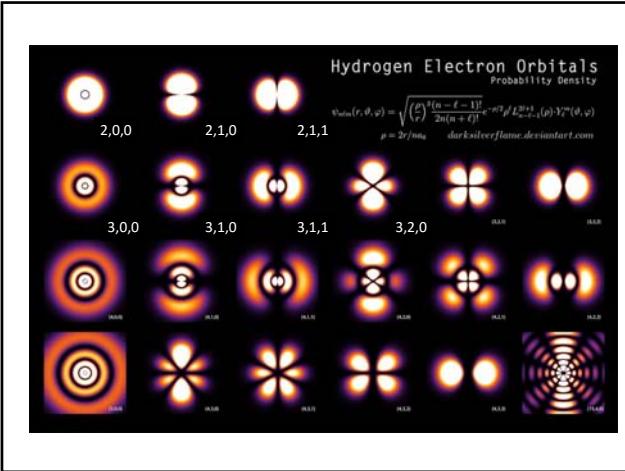
$$\Theta(\theta) = f(\theta)$$
 in our derivation in class

The first few angular functions  $\Theta_{l,m}(\theta)$ . The functions with  $m$  negative are given by  $\Theta_{l,-m} = (-1)^m \Theta_{l,m}$ .

	$l = 0$	$l = 1$	$l = 2$
$m = 0$	$\sqrt{1/4\pi}$	$\sqrt{3/4\pi} \cos \theta$	$\sqrt{5/16\pi} (3 \cos^2 \theta - 1)$
$m = 1$		$-\sqrt{3/8\pi} \sin \theta$	$-\sqrt{15/8\pi} \sin \theta \cos \theta$
$m = 2$			$\sqrt{15/32\pi} \sin^2 \theta$

$$Y_{lm}(\theta, \phi) = \Theta_{lm}(\theta) e^{im\phi}$$



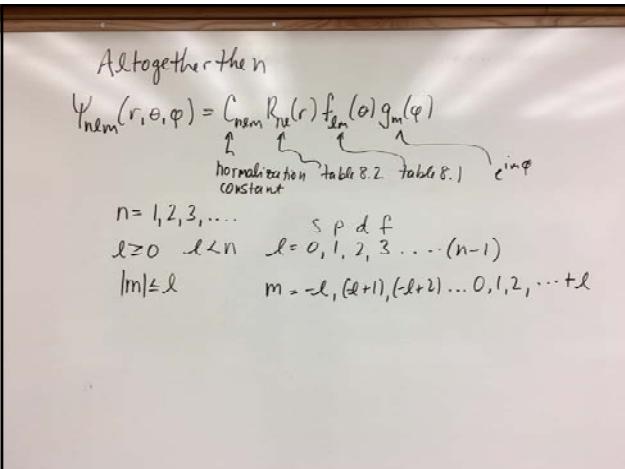


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TABLE 8.1  $\Theta(\theta) = f(\theta)$  in our derivation in class

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$\sqrt{1/4\pi}$	$\sqrt{3/4\pi} \cos \theta$	$\sqrt{5/16\pi} (3 \cos^2 \theta - 1)$
	$-\sqrt{3/8\pi} \sin \theta$	$-\sqrt{15/8\pi} \sin \theta \cos \theta$
		$\sqrt{15/32\pi} \sin^2 \theta$

$$Y_{lm}(\theta, \phi) = \Theta_{lm}(\theta) e^{im\phi}$$

Quantum number  $l$ : 0 1 2 3

Magnitude  $L$ : 0  $\sqrt{2}\hbar$   $\sqrt{6}\hbar$   $\sqrt{12}\hbar$

Code letter:  $s p d f$

$E = 0$

$E_4 = -E_R/16$   $\frac{4s}{3s}(1)$   $\frac{4p}{3p}(3)$   $\frac{4d}{3d}(5)$   $\frac{4f}{3f}(7)$

$E_3 = -E_R/9$   $\frac{3s}{3s}(1)$   $\frac{3p}{3p}(3)$

$E_2 = -E_R/4$   $\frac{2s}{2s}(1)$   $\frac{2p}{2p}(3)$

$E_1 = -E_R$   $\frac{1s}{1s}(1)$

(Energy spacing not to scale.)

code letter:  $s p d f g h i$

quantum number  $l$ : 0 1 2 3 4 5 6

$n = 1, 2, 3, \dots$

$\ell < n$

$|m| \leq \ell$

